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A NEW CLASS OF OPTICAL IMAGING SYSTEMS ACHIEVING  
'APERTURE SYNTHESIS' FROM NON-CONVENTIONAL OPTICS  
BY A POSTERIORI LENSLESS FOURIER-TRANSFORM HOLOGRAPHY \*

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The difficulties of constructing adequately perfect conventional focusing systems for some applications (e.g., X-rays and ultrasonic imaging, space applications, etc.) may be solved by a posteriori holographic synthesis.

The well-known difficulties of constructing or even of conceiving adequately perfect *conventional* focusing systems, e.g., using mirrors and lenses, for some imaging problems (e.g., X-ray astronomy) have again been recently stressed [1, 2]. For example, a solution has long been sought [1, 2] which could somehow permit one to synthesize into one single image the multiplicity of incoherently recorded images formed by such *non-conventional* systems as an aperture pierced by a large number of randomly disposed 'pin-holes' [1-4] and others [5, pp. 127-137]. Dicke [2] and Underwood [1] stressed that the required image 'synthesis' for the X-ray astronomy application could be readily achieved a posteriori, with coherent light, in an optical Fourier-transforming arrangement [5], provided that an 'appropriate phase- and amplitude-correcting plate' [2] could somehow be devised. We now

show in a general way that such a 'correcting plate' for this and a family of comparable applications may in fact be best realized in the form of a lensless Fourier-transform hologram [6] and the desired image synthesis in an 'extended-source lensless Fourier-transform' holography arrangement [5, 7]. Image 'deblurring' methods [7], at first sight formally comparable, have heretofore primarily implied 'sharpness' restoration in images imperfectly recorded with more or less 'perfect' *conventional* focusing systems, rather than being concerned with 'aperture respectively image-synthesis' of 'perfectly' recorded images, starting from suitable *non-conventional optics*, as we are here. Further analogies may be found in high-resolution X-ray crystallography [5] and radio-astronomy, among others.

Our experiments, described below, have fully born out the predicted perfection achievable with this new method of two-step imaging. Because of limitations of space, we give our theory together with the basic description of our results, used as a model for this and comparable situations. Further details will be given in ref. 4.

\* Early aspects of this work were presented on 24 September 1968 by special invitation in Florence, at the International Commission for Optics Symposium on Applications of Coherent Light.

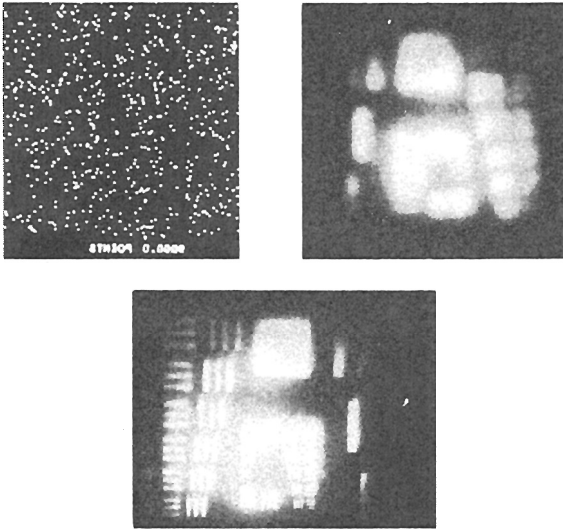


Fig. 1. A posteriori synthesis of a single image (c) from the multiple-image superposition (b) recorded with a random-array scatter-hole plate 'camera' (a) [see text]. [Note: (a) and (c) have been accidentally printed reversed with respect to (b)]. Further image improvement, to compensate for the finite 'pinhole' diameters may be achieved as in ref. 8.

Fig. 1a shows a  $4 \times 4 \text{ mm}^2$  enlarged section of the  $17 \times 17 \text{ mm}^2$  9 000 hole (0.046 mm diameter) array [3,4], having a geometrical 'projection'  $h(x, y)$  in the image plane (when illuminated by a point source at infinity) and used to form the multiple images in incoherent light. Fig. 1b shows the image

$$g(x', y') = \iint_{-\infty}^{+\infty} f(x, y) h(x' - x, y' - y) dx dy = f \otimes h$$

of a test-bar target [geometrical image =  $f(x, y)$ , about  $40 \times 40 \text{ mm}^2$  outside] which would be obtained using the pinhole array in place of a good focusing system. Because of obvious optical diffraction disadvantages resulting from diffraction at optical in comparison to X-ray radiations, as used in the optical simulation with the small pinholes, the image fig. 1b shown was actually reconstructed with point-source illumination from

the lensless Fourier-transform hologram 'equivalent' to [5]

$$I(u, v) = 1 + |\bar{G}|^2 + \bar{F}\bar{H} + \bar{F}^*\bar{H}^*$$

and obtained by using the multiple-pinhole array  $h$  as the 'extended' reference source  $h(x, y)$  in place of the point source in recording the hologram of the 'desired' image  $f(x, y)$ . We have the Fourier-transform relations

$$\bar{G}(u, v) = \iint_{-\infty}^{+\infty} g(x, y) \exp[2\pi i (ux + vy)] dx dy,$$

and similarly for  $\bar{H}$  and  $\bar{F}$ , using the usual normalization [5]. By the suitable choice of  $h \times h^* = \iint_{-\infty}^{+\infty} h(x, y) h^*(x+x', y+y') dx dy \approx \delta$  (a delta function), illumination of the hologram with light from the point-spread function  $h$  (here the projection of the aperture) used as the reconstructing source, indeed permits us to synthesize a single very good image  $f(x, y)$ , as shown in fig. 1c, from the superposition (convolution) of the multiplicity of images of fig. 1b.

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